# Molecular mechanics and molecular dynamics simulations of the hexagonal crystalline form of the helical polyamide poly( $\alpha$ -isobutyl-L-aspartate)

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The program AMBER 3.0 Rev.A has been used in a molecular mechanical study of two conformational models for the title compound on a simplified solid-state model with hexagonal coordination. Results show that the right-handed model is the minimum energy structure with a relative stability of 20 kcal mol<sup>-1</sup> residue. Molecular dynamics trajectories show less torsion angle variability in the backbone than in the side chain, according to the packing environment of each residue.

(Keywords: poly(α-isobutyl-L-aspartate); molecular mechanics; molecular dynamics)

#### INTRODUCTION

The helical structures<sup>1</sup> for the hexagonal and tetragonal<sup>2,3</sup> crystalline forms of poly( $\alpha$ -isobutyl-L-aspartate) (PAIBLA; Figure 1) were discussed in earlier papers<sup>3,4</sup>. Both X-ray diffraction analysis with the linked-atom leastsquares (LALS) methodology<sup>5</sup> and energy minimization calculations showed a right-handed helix for the tetragonal form<sup>4</sup>. On the other hand, LALS results indicate that there are two conformations, 1L and 2R (see Figure 2) (where L and R refer to a left-handed and a right-handed helix, respectively), compatible with the hexagonal form of PAIBLA4. Both models are consistent with the 13/4 helix (13 residues in four turns) observed experimentally<sup>3</sup>. In addition, theoretical calculations of the circular dichroism spectrum do not provide a distinction between them<sup>6</sup>. Preliminary energy minimization calculations favoured the model 2R<sup>4</sup> Nevertheless, these energy calculations were performed on an isolated chain, and consequently intermolecular side-chain interactions were not taken into account. In view of the importance of packing interactions in the structure of a crystalline polymer chain<sup>7-12</sup>, these results could not be considered as definitive. In the present work we have analysed the conformational and potential energies of models 1L and 2R, considering both intramolecular and intermolecular interactions simultaneously within each minimization cycle, in order to ascertain which model is energetically favoured in the hexagonal crystalline form.

Moreover, in previous LALS calculations<sup>3,4</sup> all sidechain groups were considered with the same conformation owing to the small number of diffraction spots available.

This is physically unreasonable in the hexagonal form, since depending on the packing environment the side chains could have different conformations. Thus, in order to investigate the conformational preferences of the side chains we performed a 60 ps molecular dynamics trajectory.

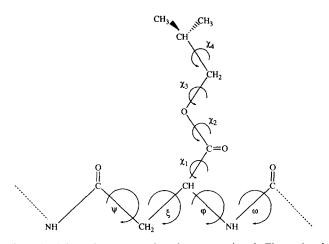


Figure 1 Schematic representation of a monomeric unit. The torsional angles are indicated. The amide group is repeated at both ends

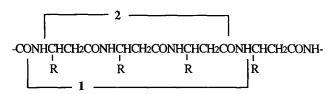


Figure 2 Different hydrogen-bonding schemes compatible with the hexagonal crystalline form conformations of  $poly(\alpha-isobutyl-L-aspartate)$ 

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#### **METHODS**

All the calculations have been performed with explicit consideration of all atoms involved 13 using the AMBER 3.0 Rev.A program<sup>14,15</sup>. Potential energy parameters obtained elsewhere4 by quantum mechanical calculations<sup>16,17</sup> were used.

Geometry optimizations were carried out in several steps. First, the starting conformations were subject to 250 cycles of steepest descent optimization to eliminate the worst steric conflicts. Second, subsequent optimization using a conjugate gradient algorithm was performed. All the structures were minimized until the difference in energy between two successive iterations was less than  $10^{-7}$  kcal mol<sup>-1</sup> (1 cal = 4.2 J) or the norm of the gradient for two successive steps in the minimization was less than  $0.1 \text{ kcal mol}^{-1} \text{ Å}^{-1} \ddagger$ .

Minimizations were carried out without any symmetry constraints. We imposed an 8 Å cut-off for the nonbonded interactions<sup>18</sup>, and updated the list of these interactions every 25 steps. Dielectric constants of  $\varepsilon = 1$ and  $\varepsilon = 1r$  provide the most reliable results in crystalline polymer systems<sup>7</sup>, and therefore both expressions were used in the energy optimization calculations.

In the molecular dynamics simulation a time step of 1 fs was used. The dielectric constant was fixed at  $\varepsilon = 1r$ . During the first 10 ps, the system was heated from 0 K to 298 K. The simulation was continued for 50 ps at 298 K. The results were averaged over the last 25 ps of the run. A configuration was stored every 1 ps for analysis.

# RESULTS AND DISCUSSION

We used the coordinates obtained from the LALS analysis performed elsewhere<sup>4</sup> to generate a polymer chain of 17 residues, which was blocked at the amino terminal end with an acetyl group and with an N-methylamide at the carboxy end. A simplified model of the hexagonal crystal form was mimicked by a central polymer chain surrounded by six equivalent chains (see Figure 3). In order to avoid end effects on the conformation, only the seven middle residues of the central chain were considered when analysing the conformation and, in addition, their dihedral angles were averaged. On the other hand, energy contributions were computed for the middle residue of the central chain1. The results obtained from the energy minimization of models 1L and 2R appear in Table 1 and Figure 3.

As can be seen, the energy-minimized conformations are not drastically sensitive to the dielectric constant considered. Thus, both expressions,  $\varepsilon = 1$  and  $\varepsilon = 1r$ , provide very similar conformations within each model. The small divergence of the minimized structures from those obtained from X-ray data refinement suggests that the original LALS conformations<sup>4</sup> are reasonably accurate and very close to an energy minimum, in particular the model 2R structure. Moreover, the peptide bond displays more reasonable values than those obtained from previous calculations on an isolated chain<sup>4</sup>. Thus, a very small deviation from planarity is observed in the peptide bond of 1L when packing

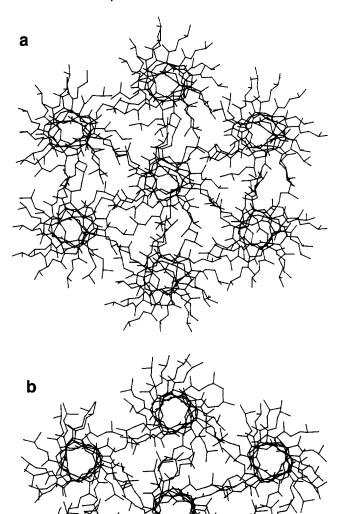


Figure 3 Minimized energy conformations of models (a) 1L and (b) **2R**, with  $\varepsilon = 1r$ , viewed along the chain axis. The different aspect of the main chain results from a better hydrogen-bonding geometry in the model 2R

interactions are taken into account, whereas a deviation of 15° is obtained when a unique chain is considered.

The total energy contributions indicate that model 2R is approximately 20 kcal mol<sup>-1</sup> residue more favourable with respect to model 1L. Electrostatic, van der Waals and bonding energetic contributions favoured the model 2R by around 6 kcal mol<sup>-1</sup> residue each, for both dielectric constants. Model 2R preserves better the average hydrogen-bonding geometry than model 1L, as we see in Figure 3: the helices in model 1L clearly appear more distorted at the main chain. Furthermore, the similarities between the model 2R of the hexagonal form and the right-handed helix of the tetragonal form allow us to explain some experimental observations, e.g. that the transition between the hexagonal and tetragonal forms takes place easily<sup>2</sup>, and that both forms have very similar n.m.r. spectra in the solid state<sup>19</sup>.

<sup>‡</sup> The different energy terms refer to the middle residue of the central polymer chain, where both interresidue and intraresidue interactions were taken into account

Results obtained from the energy minimization of models 1L and 2R using the two different expressions for the dielectric constant

	Disloctrio				Conformational angles (°)	nal angles	(,)			Dielectric			Calculated	Calculated energies (kcal mol <sup>-1</sup> residue)	mol <sup>-1</sup> resic	ne)
iel"	constant	ф	ν,	<b>/</b>	Ø	χ1	χ2	χ₃	χ4	constant	Contribution	$E_{ m bond}$	Evdw	$E_{ m el}$	$E_{ m hb}$	$E_{ m tot}$
	1	149.5	-93.0	9.68	175.0	179.7	-170.4	-71.8	138.5	1	Intra	14.88	7.37	-124.63	-0.42	-102.81
	11	161.7	-96.8		-177.2	174.4	185.6	- 68.8	144.7		Inter	3.56	-3.53	-73.65	-1.07	-101.68
	$LALS^b$	126.6	-102.8	113.6	180.0	175.0	180.0	-101.2	114.1		Total	18.44	-23.16	-198.29	-1.49	- 204.49
										11	Intra	14.76	7.02	-36.40	-0.09	-14.70
											Inter	4.82	-31.81	-38.95	-0.53	-66.47
											Total	19.58	-24.79	-75.35	-0.63	-81.17
	-	157.2	-54.1	114.8	175.0	173.4	172.4	-108.1	112.7	1	Intra	10.99	6.54	-111.48	-0.48	-94.43
	11	149.6	-54.0	120.8	177.9	177.2	169.0	-101.1	107.2		Inter	1.49	-36.45	-94.02	-1.19	-130.17
	LALS	145.2	-62.0	130.8	180.0	175.0	180.0	-101.2	114.1		Total	12.48	-29.91	-205.50	-1.67	-224.60
										11	Intra	11.37	4.63	-29.63	-0.46	- 14.09
											Inter	0.80	-35.12	-52.32	-0.76	-87.40
											Total	12.17	-30.48	-81.96	-1.23	-101.50

**2R** refer to a left-handed and a right-handed helix, respectively, with the hydrogen-bonding schemes 1 and 2 shown in Figure 2. The conformational angles correspond to those indicated in The values given are the averages of the middle seven residues of the central chain LALS-refined<sup>4</sup> values are shown for comparisor Models 1L and 2R refer igure 2 of ref.

to one residue, where  $E_{bord} = bonding$  energy,  $E_{cdw} = van$  der Waals energy,  $E_{cl} = electrostatic$  energy,  $E_{bb} = energy$  corresponding to the  $r^{-12} - r^{-10}$  term and  $E_{cd} = total$  energy. Let an an intrahelical and interhelical interactions. The starting conformations were calculated with the LALS program taking <sup>c</sup>The different energy terms refer to one Inter and intra refer to interresidue and into account the X-ray results<sup>4</sup>

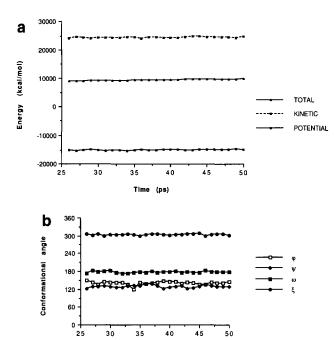


Figure 4 (a) Energy and (b) backbone conformational angle fluctuations along the last 25 ps of the molecular dynamics trajectory

**Table 2** Averaged side-chain conformational angles and standard deviations (in parentheses) for the seven central residues along the last 25 ps of the trajectory. The total average values are also displayed for comparison with the LALS values (see *Table 1*)

	Conformational angles (°)				
Residue	χ <sub>1</sub>	χ <sub>2</sub>	χ <sub>3</sub>	χ4	
1	168.1	-177.7	-169.8	58.7	
	(14.7)	(11.7)	(16.6)	(8.7)	
2	$-169.8^{\circ}$	-173.8	-158.1	178.2	
	(33.6)	(17.7)	(39.1)	(11.0)	
3	134.2	-172.6	<b>–</b> 75.9	$-179.2^{\circ}$	
	(14.0)	(15.0)	(44.0)	(9.0)	
4	-136.1	$-165.0^{\circ}$	174.4	$-64.8^{'}$	
	(16.3)	(14.3)	(19.8)	(18.7)	
5	$-156.9^{'}$	- ì71.9 <sup>°</sup>	$-172.3^{\circ}$	$-62.3^{'}$	
	(19.8)	(11.1)	(20.3)	(20.4)	
6	-132.1	ì71.5 <sup>°</sup>	ì76.7	54.5	
	(12.5)	(12.9)	(12.3)	(10.1)	
7	177.2	179.4	-173.2	59.6	
	(20.3)	(12.7)	(10.2)	(12.0)	
Total average	-170.8	-173.3	-159.7	160.7	

The energy-minimized coordinates of model 2R were taken as the starting point of a molecular dynamics trajectory. Figure 4 shows the fluctuations of the total energy of the system and the backbone torsion angles. The simulation reproduces reasonably well the refined X-ray structure. There are no large fluctuations in the backbone torsion angles, indicating small deviations relative to the X-ray results. These represent a very satisfactory agreement considering that energetic simulations have been carried out without geometrical constraints on the atomic coordinates.

Table 2 displays the side-chain conformational angles and standard deviations along the molecular dynamics trajectory for the seven central residues. An overall picture shows that whereas backbone conformational angles are conserved at the same value for each residue

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along the trajectory, the side-chain dihedral angles display some variations for each residue depending on the packing environment. The seven central residues show different conformational angles and, in general, standard deviations are small. The lowest standard deviations correspond as might be expected to the ester group  $(\chi_2)$ , which has a conformational angle of around 180°. The largest standard deviations computed along the trajectory indicate that  $\chi_3$  is the most flexible angle, whereas  $\chi_1$ and  $\chi_4$  display similar values. Both side-chain angles  $\chi_1$  and  $\chi_3$  display an extended conformation, with the exception of residue 3 for which  $\chi_3$  corresponds to a gauche conformation. On the other hand, the results indicate a gauche (residues 1, 4, 5, 6 and 7) or a trans (residues 2 and 3) conformation for the  $\chi_4$  angle, depending on the interactions with the neighbouring chains. Thus, it seems that the packing environment affects considerably the conformation at the end of the side chain. The totally average conformational angles show a close resemblance to the LALS values, indicating that although they are physically unreasonable, they provide a good average picture of the side-chain conformation.

In summary, model **2R** is the minimum energy structure for the hexagonal form of PAIBLA by around 20 kcal mol<sup>-1</sup> residue. This energy gap between models **2R** and **1L** is conserved for the different expressions of the dielectric constant. On the other hand, the molecular dynamics simulation shows that side-chain groups have different conformations depending on the packing environment generated by the neighbouring chains. Although these results could not be considered quantitatively owing to the limited duration of the simulation and the packing approximations, they provide a qualitative picture of the situation of side-chain conformations in the hexagonal crystalline form.

The present study represents an improvement of our previous work<sup>4</sup>. The results suggest that the combination of force-field calculations and refined crystallographic data seems to be very promising in structure determination, particularly in polymer fibres where poor crystallographic data are obtained.

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